

Metales y sistemas magnéticos

Plan

Metales a temperatura finita

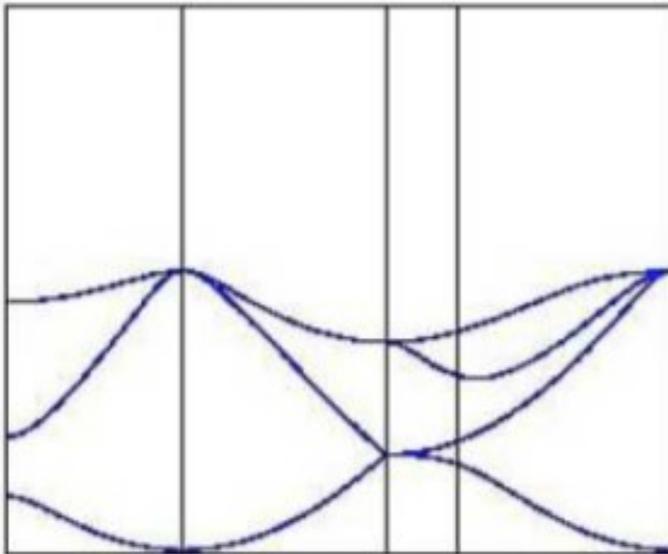
Metales a 0K. Metodo de smearing.

Sistemas magnéticos.

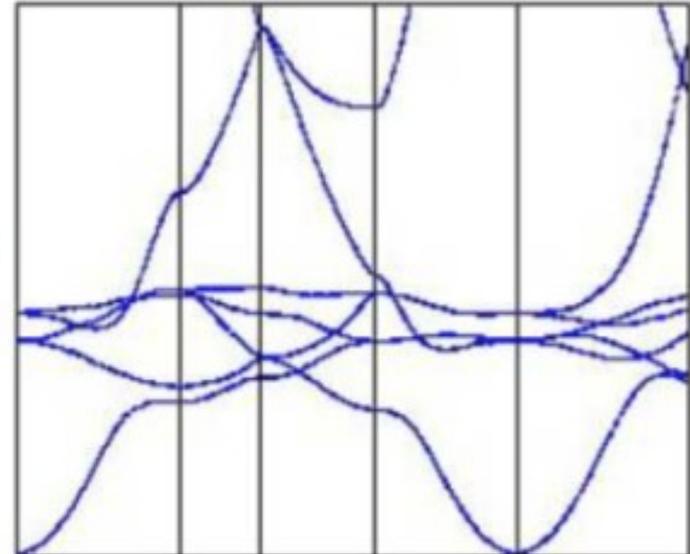
Brillouin Zone Integration

1. Sampling at one point (Gamma or Baldereschi point)
2. Sampling at regular grids (Monkhorst-Pack)
3. For metallic systems, integration of the discontinuity is improved introducing a fictitious electronic temperature

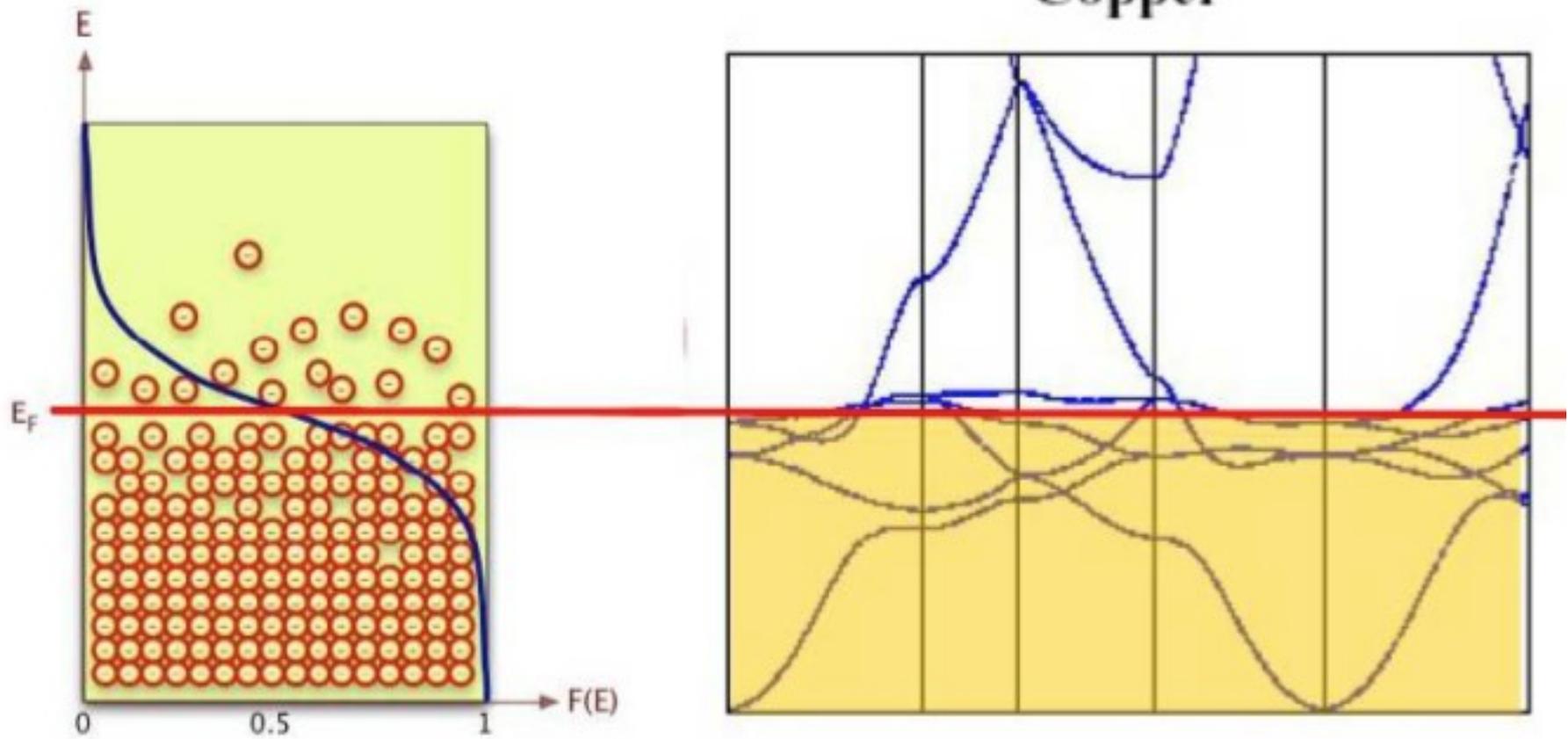
Valence bands



Copper



Copper



$$f_{n,\mathbf{k}} = \frac{1}{1 + \exp(\beta(\epsilon_{n,\mathbf{k}} - \epsilon_F))}$$

Temperatura

$$f_i = \frac{1}{\exp[(\varepsilon_i - \mu)/k_B T] + 1}$$

$$n(\vec{r}) = \sum_i^{\infty} f_i |\psi_i(\vec{r})|^2$$

$$E(T) = \sum_i f_i \varepsilon_i - \frac{1}{2} \int n(\vec{r}) V_H(\vec{r}) d^3 \vec{r} + E_{xc}[n] - \int n(\vec{r}) v_{xc}(\vec{r}) d^3 \vec{r}.$$

$$F = E - TS$$

$$S = - \sum_i [f_i \ln f_i + (1 - f_i) \ln(1 - f_i)].$$

Necesidad del smearing en metales (T= 0K)

$$\sum_n \frac{1}{\Omega_{BZ}} \int_{\Omega_{BZ}} \epsilon_{nk} \Theta(\epsilon_{nk} - \mu) dk, \longrightarrow \sum_k w_k \epsilon_{nk} \Theta(\epsilon_{nk} - \mu),$$

La suma converge lentamente con el numero de k-points en los metales. Esto ocurre porque para los k cercanos a la superficie de Fermi, los numeros de ocupacion saltan entre 0 y 1. El truco es remplazar la funcion de salto por una función mas suave, que haga suaves los cambios de ocupación al variar los niveles de energía.

$$f\left(\frac{\epsilon - \mu}{\sigma}\right) = \frac{1}{\exp\left(\frac{\epsilon - \mu}{\sigma}\right) + 1}$$

Fermi smearing

$$f\left(\frac{\epsilon - \mu}{\sigma}\right) = \frac{1}{2} \left(1 - \operatorname{erf} \left[\frac{\epsilon - \mu}{\sigma} \right] \right)$$

Gaussian smearing

Smearing ($T=0K$)



En metales se necesitan muchos \vec{k} para que converja la energía.

A veces no converge el ciclo SCF.

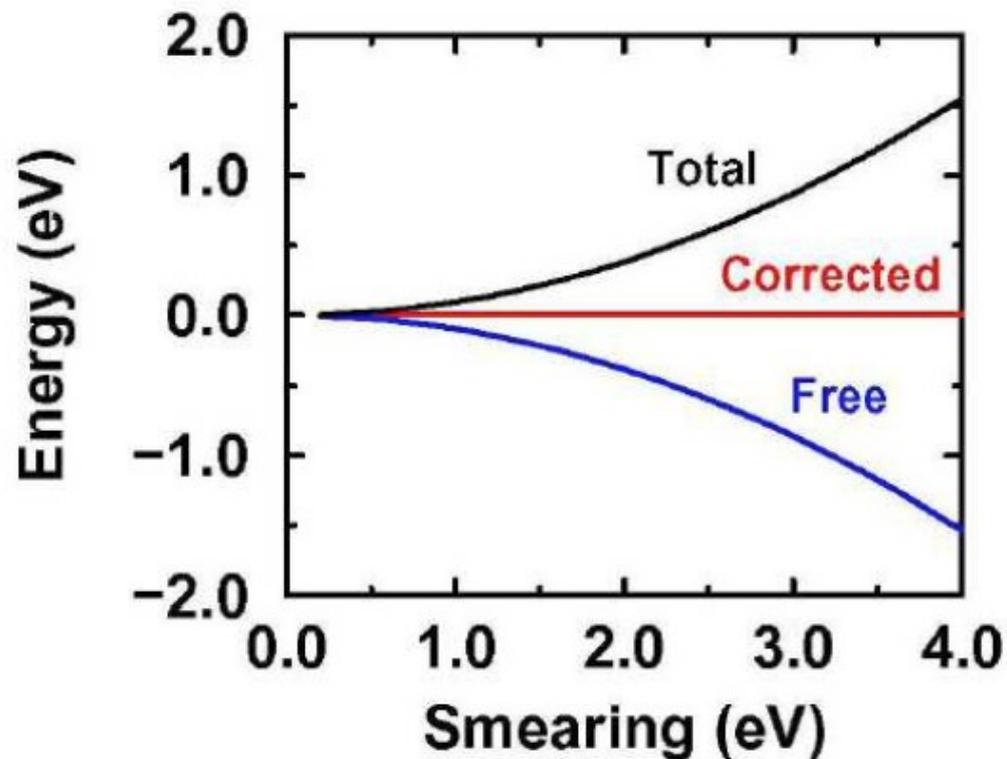
Se usa una temperatura ficticia ($k_B T \sim 0,2 \text{ eV}$) para reducir el número de puntos k . $k_B \times 300K \sim 0,025 \text{ eV}$.

$$E(T) = E(0) + \alpha T^2$$

$$F(T) = E(0) - \alpha T^2$$

$$E(0) = \frac{1}{2}(E(T) + F(T)) + O(T^4).$$

$$A[\sigma; \{\psi_i\}, \{f_i\}] = \sum_i f_i \langle \psi_i | \hat{T}_e + \hat{V}_{nl} | \psi_i \rangle + E_{\text{Hxc}}[n] - \sigma S[\{f_i\}].$$



Temperatura

Las fuerzas de Hellman-Feynman corresponden a $F(T)$, no a $E(0)$.

$$\vec{F} = -\nabla F(T)$$

Metales: $TS < 1 \text{ meV/atomo}$

Ahorro CPU, RAM \Rightarrow Precision

Pocos $k \Rightarrow$ Baja T

Alta $T \Rightarrow$ Muchos k

Metales

- Fermi smearing es ineficiente. Solo se usa para calcular energía libre verdadera.
- Otros métodos: gaussiano, Methfessel-Paxton, cold smearing.
`smearing='gaussian', 'mp', degauss=0.02`
- Reducen TS y permiten usar menos k-points.
Phys. Rev. B 65, 035111 (2001), Smearing scheme finite-temperature electronic calculations

El método de tetrahedros es una interpolación de las energías en puntos intermedios que no están en la malla de puntos k. Es lo mejor para calcular energía.

Smearing

in the **&SYSTEM** namelist

occupation CHARACTER

'smearing': smearing for metals

requires a value for degauss

'tetrahedra': for metals and DOS calculations

(see PRB 49, 16223 (1994))

requires uniform grid of k-points,
automatically generated

'fixed': for insulator with gap

'from_input': the occupation are read from input file.

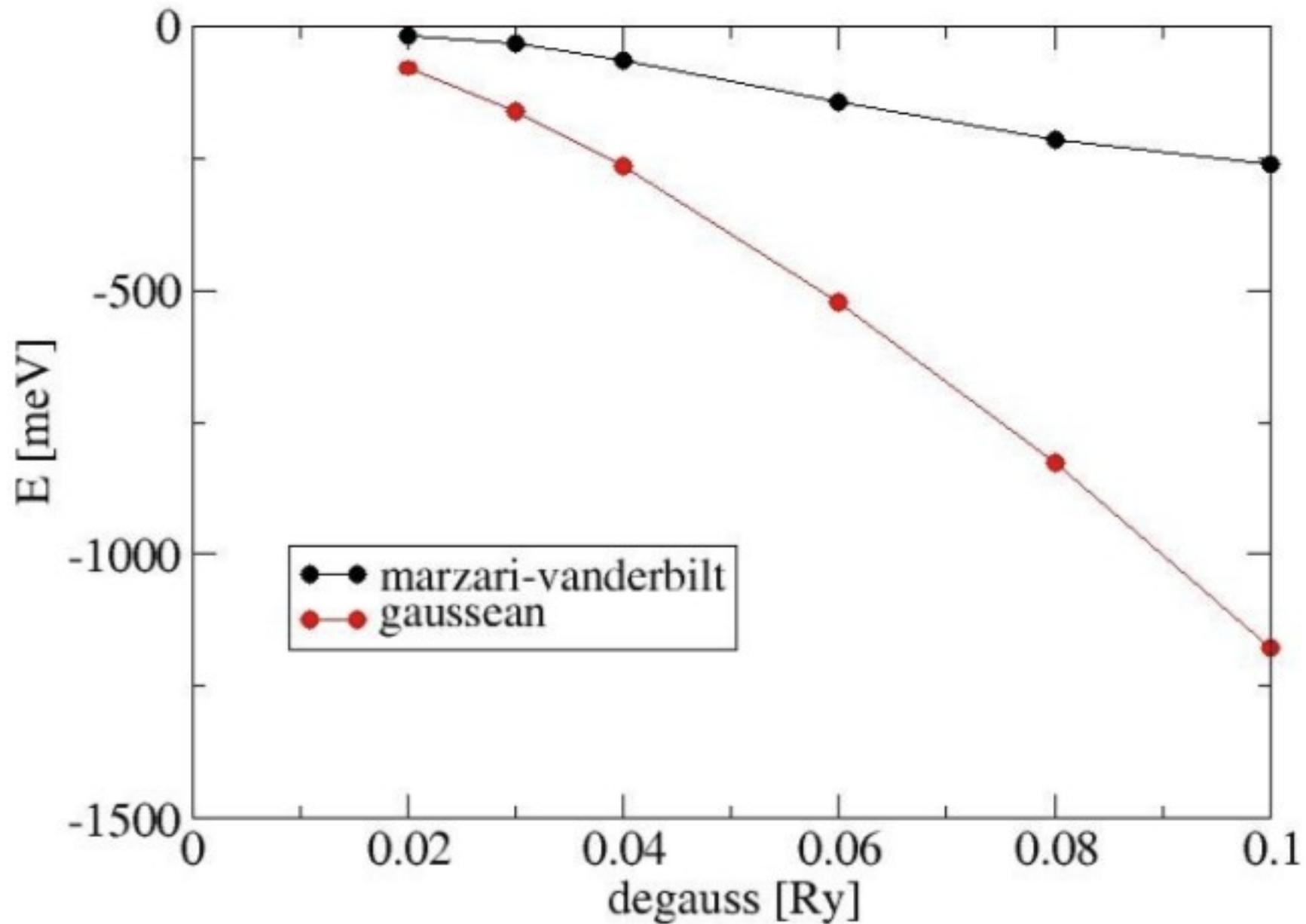
Presently works only with one-kpoint

(LSDA allowed)

**El método de tetrahedros es bueno para la energía,
mas no da las fuerzas correctas. Si necesita fuerzas,
use occupation='smearing', smearing='mp',degauss=?**

degauss REAL (default = 0.d0 Ry)
value of the smearing for brillouin-zone
integration in metals.

smearing CHARACTER
'gaussian', 'gauss' :
 ordinary gaussian smearing (default)
'methfessel-paxton', 'm-p', 'mp':
 Methfessel-Paxton first order smearing
 (see PRB 40, 3616 (1989))
'marzari-vanderbilt', 'cold', 'm-v', 'mv':
 Marzari-Vanderbilt cold smearing
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'fermi-dirac', 'f-d', 'fd':
 smearing with Fermi-Dirac function



sistemas magnéticos

So far in this course: we “neglected” spin of the electrons:

$$n(\mathbf{r}) = 2 \sum_{i=1}^{N_{el}/2} |\psi_i(\mathbf{r})|^2$$

Implicitly, we have assumed that spin-up and spin-down wavefunctions are the same:

$$\psi_i^{\uparrow}(\mathbf{r}) = \psi_i^{\downarrow}(\mathbf{r})$$

Advantage: we need only $N_{el} / 2$ orbitals

sistemas magnéticos

For the case of spin-polarized systems, we drop this constraint and Use N orbitals:

$$\left\{ \psi_i^\uparrow \right\}, \quad i = 1, \frac{N}{2} \quad \left\{ \psi_i^\downarrow \right\}, \quad i = 1, \frac{N}{2}$$

$$n^\uparrow(\mathbf{r}) = \sum_i \left| \psi_i^\uparrow(\mathbf{r}) \right|^2 \quad n^\downarrow(\mathbf{r}) = \sum_i \left| \psi_i^\downarrow(\mathbf{r}) \right|^2$$

$$n(\mathbf{r}) = n^\uparrow + n^\downarrow$$

$$m(\mathbf{r}) = n^\uparrow - n^\downarrow$$

sistemas magnéticos

Look at how the total energy depends on the two charge densities:

$$E[n^\uparrow, n^\downarrow] = E_{kin} + E_{ext}[n^\uparrow + n^\downarrow] + E_{Har}[n^\uparrow + n^\downarrow] + E_{xc}[n^\uparrow, n^\downarrow]$$

Calculate the functional derivative in order to obtain the Hamiltonian:

$$\frac{\delta E}{\delta n^\uparrow(\mathbf{r})} \neq \frac{\delta E}{\delta n^\downarrow(\mathbf{r})}$$

sistemas magnéticos

We can also write the exchange–correlation energy as $E_{xc}[n, m]$.

Since we have:

$$\frac{\delta n}{\delta n^\uparrow(\mathbf{r})} = \frac{\delta n}{\delta n^\downarrow(\mathbf{r})} = 1$$
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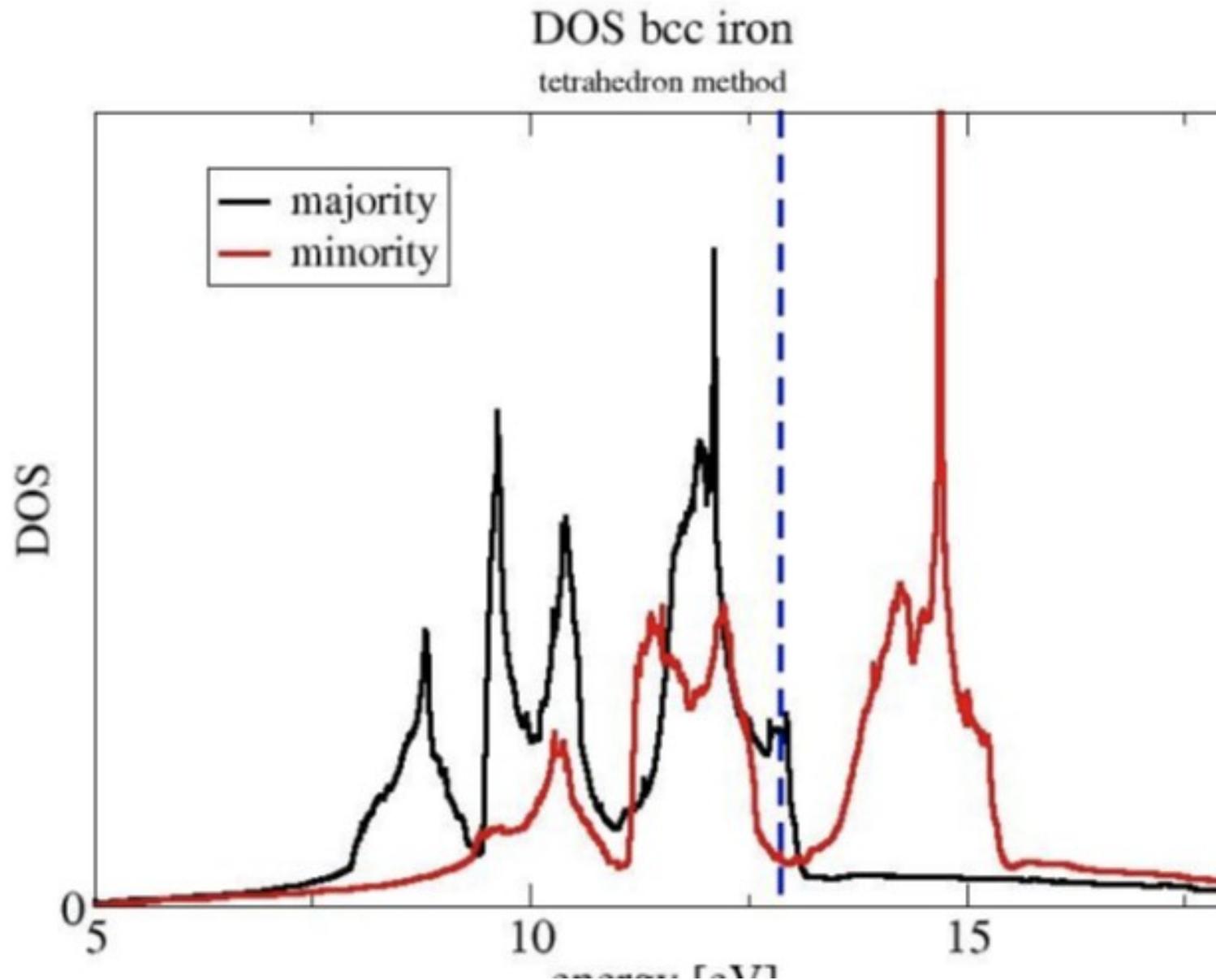
The Hamiltonian looks like:

$$\left[\frac{-\hbar^2}{2m} \nabla^2 + V_{ext} + V_{Har} + V_{xc} + B_{xc} \right] \psi_i^\uparrow = \epsilon_i \psi_i^\uparrow$$
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sistemas magnéticos

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Densidad de estados



Densidad de estados

$$D(E) = \sum_{n, \mathbf{k}, \uparrow, \downarrow} \delta(E - \epsilon_{n, \mathbf{k}})$$

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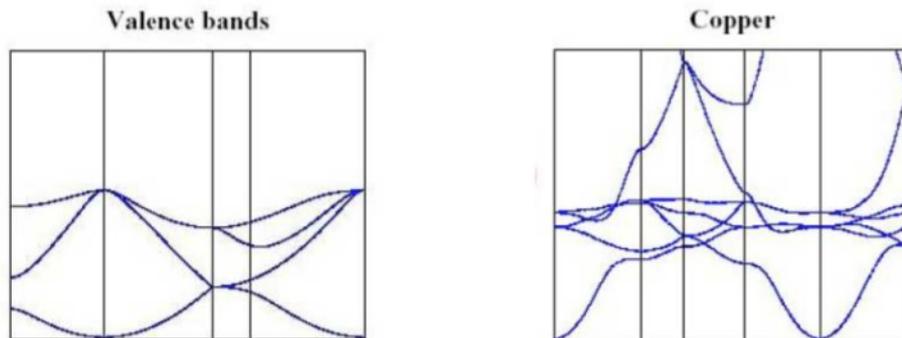
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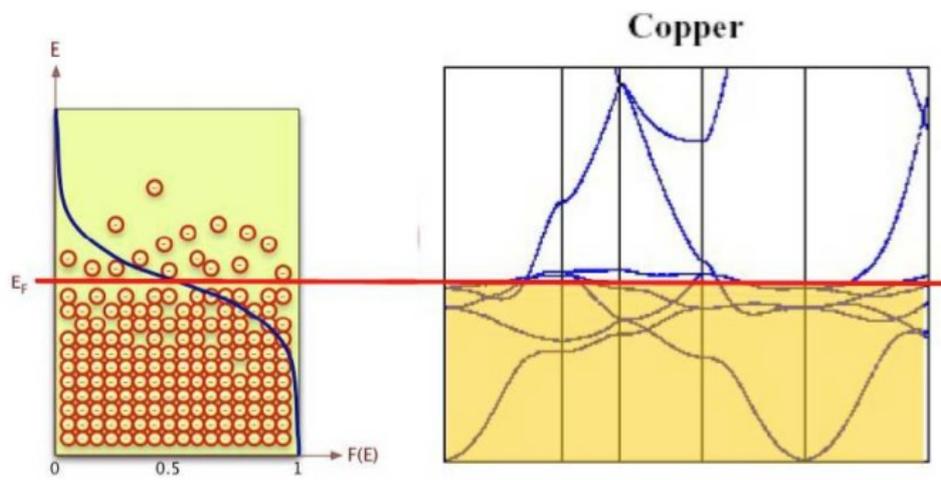
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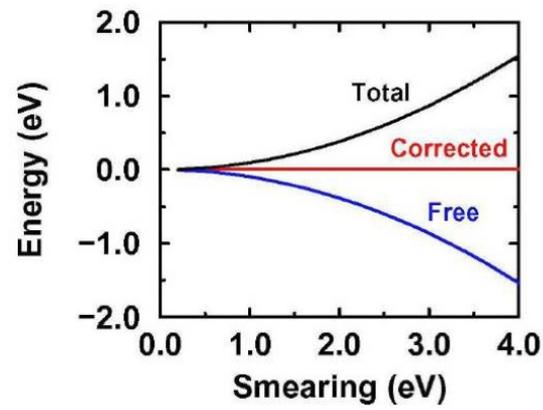
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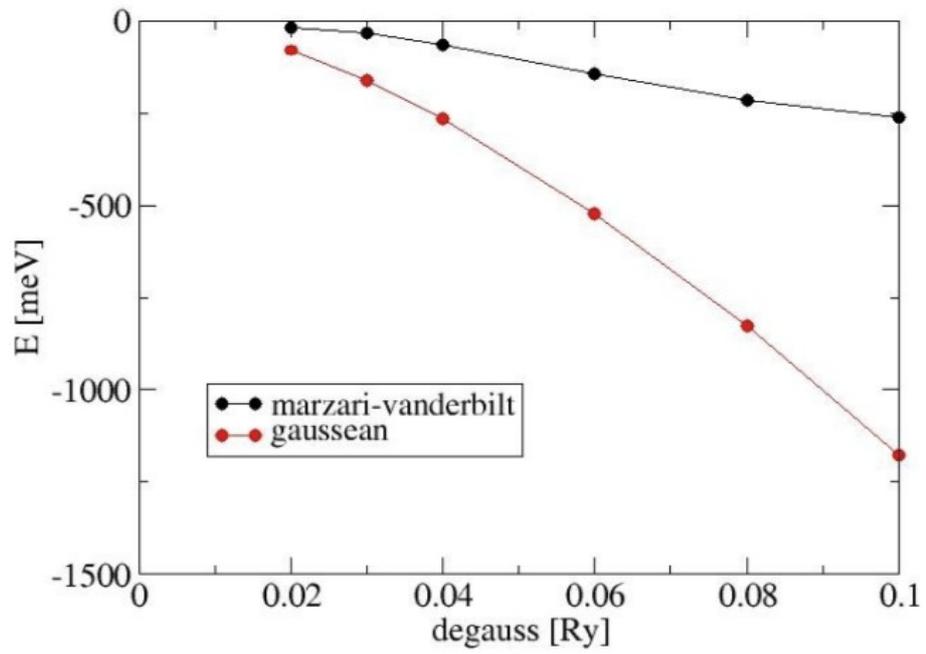
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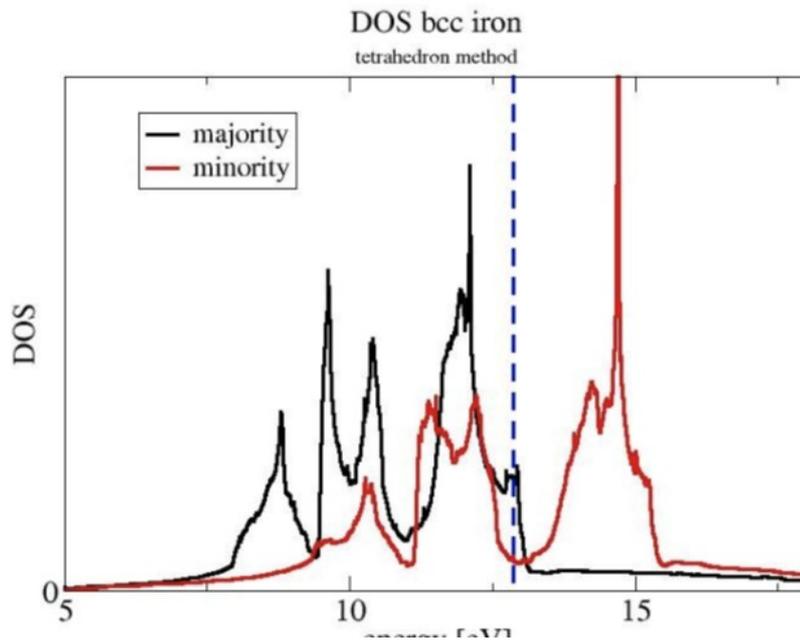
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sistemas magnéticos

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$\frac{\delta E_{xc}[n, m]}{\delta n(\mathbf{r})}$ $\frac{\delta E_{xc}[n, m]}{\delta m(\mathbf{r})}$

Densidad de estados



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